

Preparing Thermal States of Quantum Systems by Dimension Reduction

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Joint work with Sergio Boixo

Outline

- 1 Introduction
 - Motivation
 - Main Results
- 2 The Algorithm
 - Overview
 - How it works
- 3 Summary

Motivation

- **Very few quantum systems have analytical solutions.**
- Have to resort to numerical simulations in many cases
 - Brute force calculations take $\mathcal{O}(e^N)$ time and memory for N-particle systems.
 - Classical algorithms to approximate solutions (DMRG, PEPS, BP, etc) only work for specific cases.
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Previous work

- Feynman suggested quantum machines to simulate quantum systems.
- Quantum computers are very good at simulating unitary evolutions (Lloyd).
- Initial state preparation is still a difficult problem.
- Several Proposals:
 - Evolving with a bath (Terhal and DiVincenzo)
 - Quantum Metropolis Sampling (Temme *et al.*, Yung and Aspuru-Guzik)
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1D systems	$\mathcal{O}(e^{\alpha N})$	$\mathcal{O}(N^{\beta \ h\ })$
D-dimensions	$\mathcal{O}(\exp(\alpha N^D))$	$\mathcal{O}(\exp(N^{D-1}))$

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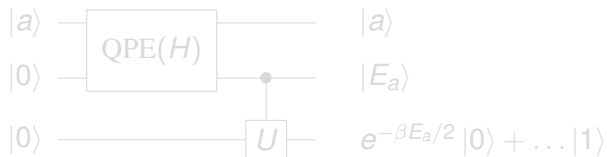
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Thermalization using QPE

- Given $H = \sum_a E_a |a\rangle\langle a|$, we want $\rho \propto \sum_a e^{-\beta E_a} |a\rangle\langle a|$.



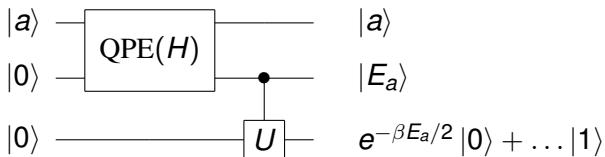
- Now, instead of $|a\rangle$, we input $I = \sum_a |a\rangle\langle a|$

$$\rightarrow \sum_a e^{-\beta E_a} |a\rangle\langle a| \otimes |E_a\rangle\langle E_a| \otimes |0\rangle\langle 0| + \dots$$

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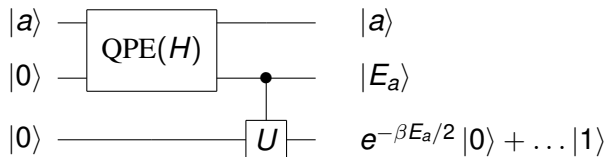
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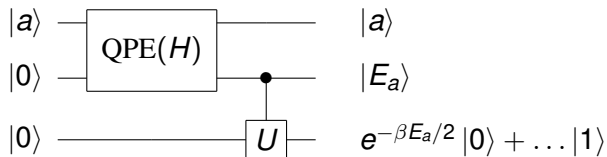
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- Projecting everything in one step costs $\mathcal{O}(e^{\beta\|H\|}) \sim \mathcal{O}(e^N)$.
- We want to break up the projections so that only a small section needs to be restarted after a failure.
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$$p \sim e^{-\beta\|h\|} \quad , \quad \text{Total cost: } \mathcal{O}(N^{\beta\|h\|})$$

Perturbative Hamiltonian Update

- We need the map $e^{-\beta H} \rightarrow e^{-\beta(H+h)}$.
- Defining $\rho^{(\epsilon)} \propto e^{-\beta(H+\epsilon h)}$, we want the sequence:

$$\rho^{(0)} \rightarrow \rho^{(\epsilon)} \rightarrow \rho^{(2\epsilon)} \rightarrow \dots \rightarrow \rho^{(1)}$$

- Each step is correct up to an error of $\mathcal{O}(\epsilon^2)$, resulting an overall error of $\mathcal{O}(\epsilon)$.

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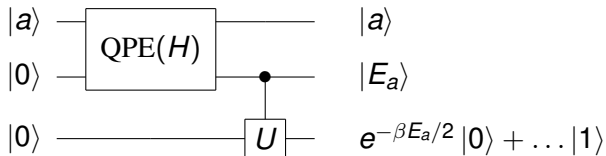
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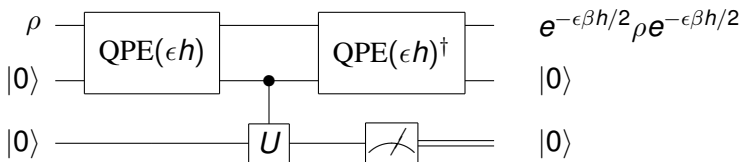
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with probability $p \geq e^{-\epsilon\beta \|h\|}$.

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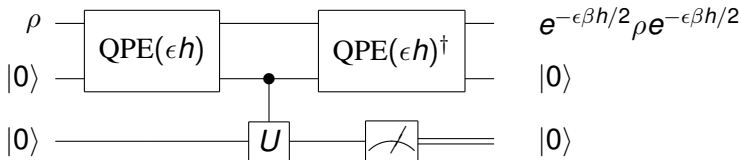
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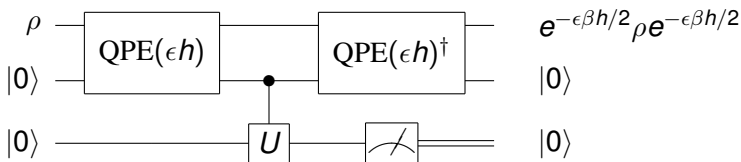
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- We dephase in the eigenbasis of the new Hamiltonian, $H + \epsilon h$.
- After the QPE circuit, we had $\rho_{\text{prob}} \propto e^{-\epsilon\beta h/2} \rho e^{-\epsilon\beta h/2}$.
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Putting Everything Together

- We can now implement the map $e^{-\beta H} \rightarrow e^{-\beta(H+h)}$ using the sequence:

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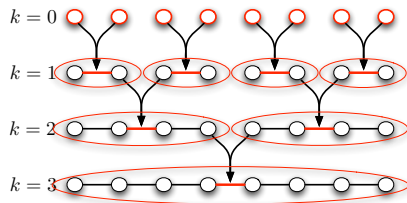
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- To thermalize a chain of $N = 2^k$ qubits, we use the recursive merging procedure from earlier:

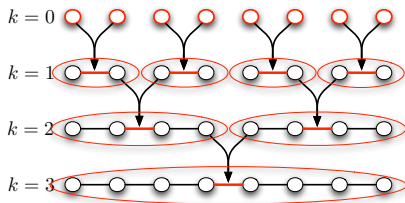


$$\tau(k) = \alpha 2 \tau(k-1) + m$$

- For an error $\bar{\epsilon}$, running time for 1D: $\tau \sim \beta N^{\beta \|h\|} / \bar{\epsilon}^2$
- For D-dimensions: $\tau \sim \beta e^{2\beta \|h\|} D N^{D-1} / \bar{\epsilon}^2$

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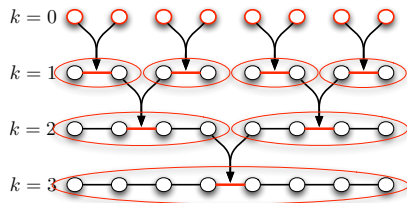


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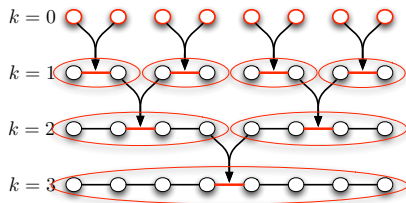


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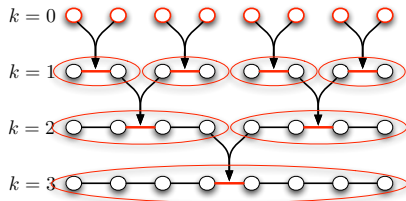
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- Made possible by recursively merging smaller regions using QPE and dephasing



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