

Large violation of Bell inequalities with low entanglement

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LHVM vs Quantum Mechanics

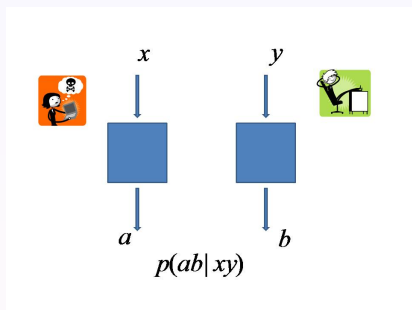


Figure: Alice and Bob measurements. Inputs: x and y , Outputs: a and b . $P(a, b|x, y)$ is the probability of obtaining the pair (a, b) when Alice and Bob measure, respectively, with the input x and y .

We deal with $(P(a, b|x, y))_{x,y=1,\dots,N}^{a,b=1,\dots,K} \in \mathbb{R}^J (J = N^2 K^2)$.

Probability distributions

Classical probabilities:

$$P = P(a, b|x, y) = \int_{\Omega} P_{\omega}(a|x)Q_{\omega}(b|y)d\mathbb{P}(\omega),$$

- a) (Ω, P) is a probability space,
- b) $P_{\omega}(a|x) \in \{0, 1\}$ and $\sum_{a=1}^K P_{\omega}(a|x) = 1$ for every x, a, ω .

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Quantum probabilities: $P(a, b|x, y) = \text{tr}(E_x^a \otimes F_y^b \rho)$,

- a) ρ is a density operator acting on $H_1 \otimes H_2$.
- b) $E_x^a \geq 0$ for every x, a and $\sum_{a=1}^K E_x^a = \mathbb{1}$ for every x (and analogously for F_y^b).

$$\mathcal{Q} = \{P : P \text{ is quantum}\}$$

“Distance” between quantum and classical probability distributions

Given $M = \{M_{x,y}^{a,b}\}_{x,y=1,a,b=1}^{N,K}$, we define

$$LV(M) = \frac{\sup\{|\langle M, Q \rangle| : Q \in \mathcal{Q}\}}{\sup\{|\langle M, P \rangle| : P \in \mathcal{L}\}} = \frac{\omega^*(M)}{\omega(M)}.$$

Here,

$$\langle M, P \rangle = \sum_{x,y,a,b=1}^{N,K} M_{x,y}^{a,b} P(a, b|x, y).$$

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Note that

$$d(\mathcal{L}, \mathcal{Q}) = f(N, K, d).$$

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- 1- This is what we need to study $\omega^*(M)$
 \rightsquigarrow Difference between *operator algebras* and *operator spaces*
- 2- Equivalently:

$$M_n \otimes E, \quad n \geq 1$$

$LV(M) \rightsquigarrow$ classical theory vs non-commutative theory

A nice construction

Let $n \in \mathbb{N}$. Consider $\epsilon_{x,a}^k = \pm 1$ with $x, a, k = 1, \dots, n$ and define

$$u_x^a = (1, \epsilon_{x,a}^1, \dots, \epsilon_{x,a}^n) \in \mathbb{R}^{n+1} \quad \text{for every } x, a = 1, \dots, n.$$

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a) *Bell inequality coefficients:*

$$M_{x,y}^{a,b} = \begin{cases} \frac{1}{n^2} (\langle u_x^a, u_y^b \rangle - 1) & x, y, a, b = 1, \dots, n \\ 0 & a = n+1 \text{ or } b = n+1. \end{cases}$$

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b) *POVMs measurements:* $\{E_x^a\}_{x,a=1}^{n,n+1}$ in M_{n+1} : For $x = 1, \dots, n$

$$E_x^a = \begin{cases} |\tilde{u}_x^a\rangle\langle\tilde{u}_x^a| & \text{for } a = 1, \dots, n, \\ 1 - \sum_{a=1}^n E_x^a & \text{for } a = n+1. \end{cases}$$

Here $\tilde{u}_x^a = \frac{1}{\sqrt{nK}} u_x^a$ for certain universal constant K .

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- 1) $\sup\{|\langle M, P \rangle| : P \in \mathcal{L}\} \leq \log n$.
- 2) For any (diagonal) pure state $|\psi\rangle = \sum_{i=1}^{n+1} \alpha_i |ii\rangle$ we have

$$|\langle M, Q_{|\psi}\rangle| \geq \alpha_1 \left(\sum_{i=2}^{n+1} \alpha_i \right),$$

where $Q_{|\psi}(a, b|x, y) = \langle \psi | E_x^a \otimes E_y^b | \psi \rangle$, $x, y, a, b = 1, \dots, n$.

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Taking $|\varphi_\alpha\rangle = \alpha |11\rangle + \frac{\sqrt{1-\alpha^2}}{\sqrt{n}} \sum_{i=2}^{n+1} |ii\rangle \in \ell_2^{n+1} \otimes \ell_2^{n+1}$ we have

$$LV(M) \geq \alpha \sqrt{1-\alpha^2} \frac{\sqrt{n}}{\log n}.$$

Consequence I: Large violation of Bell Inequalities

Theorem

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In the previous talk the authors showed

Theorem (Buhrman, Regev, Scarpa, de Wolf)

$$f\left(\frac{2^n}{n}, n, n\right) \succeq \frac{n}{(\log n)^2}.$$

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Theorem

For any $\delta > 0$ we can find a n -dimensional pure state $|\psi_\delta\rangle$ in a high enough dimension n verifying:

a) $\mathcal{E}(|\varphi\rangle) < \delta$ (resp. $\log_2(n) - \mathcal{E}(|\varphi\rangle) < \delta$),

b)

$$\frac{|\langle M, Q_{|\psi_\delta\rangle} \rangle|}{\sup_{P \in \mathcal{L}} |\langle M, P \rangle|} \succeq \frac{\sqrt{n}}{(\log n)^2},$$

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Even though quantum entanglement is needed to obtain violation of Bell inequalities, the amount of entanglement is essentially irrelevant for large violations.

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There is a Bell inequalities M with 2^{n^2} inputs and n outputs s.t.

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$$LV_n(M) = \frac{\sup\{|\langle M, Q \rangle| : Q \in \mathcal{Q}_n\}}{\sup\{|\langle M, P \rangle| : P \in \mathcal{L}\}} \succeq \frac{\sqrt{n}}{\log n},$$

b) $\sup\{|\langle M, Q \rangle| : Q \in \mathcal{Q}_{max}\} \leq 1$, where \mathcal{Q}_{max} is the set of quantum probability distributions constructed with the maximally entangled state in *any* dimension.

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- 1) There exist quantum probability distributions which cannot be written by using *any* maximally entangled state.
- 2) “Opposite” result to the main one in the previous talk.

THANK YOU VERY MUCH FOR YOUR ATTENTION